

# Hydrodynamic Limits in Particle Systems

Julian Kern  
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# Disclaimer

- not rigorous at all!
- mostly heuristics
- references at the end

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- slides for experts...

## 1 The Symmetric Simple Exclusion Process

## 2 What Are Hydrodynamic Limits

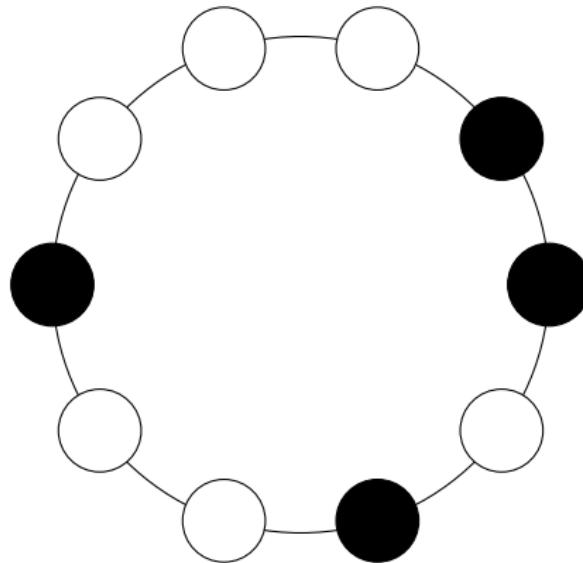
## 3 Extending the Model

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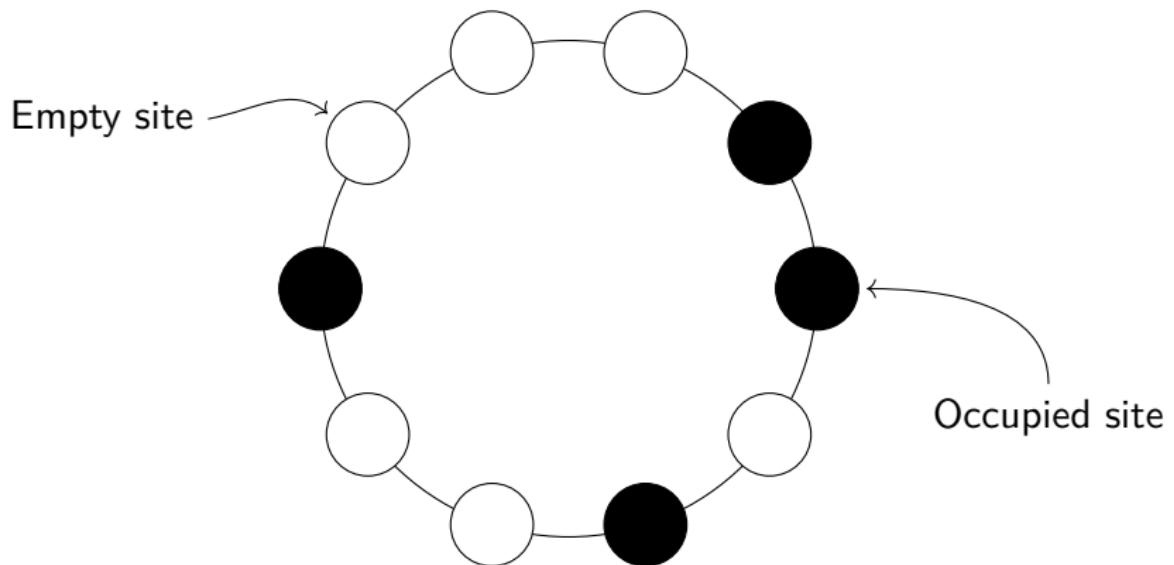
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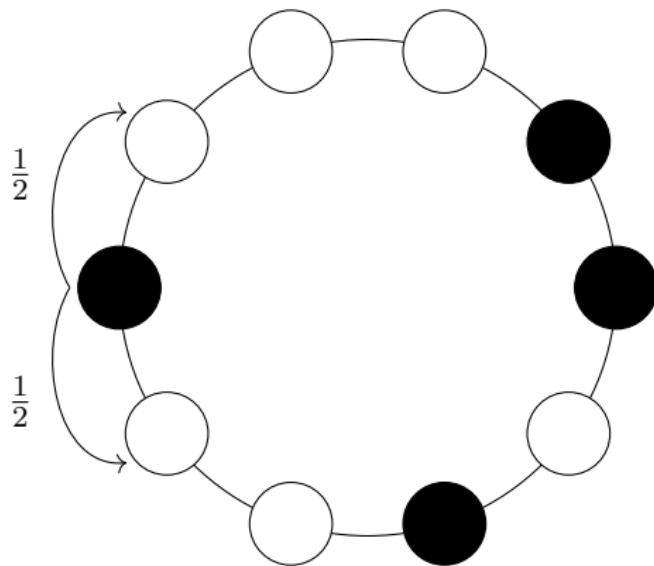
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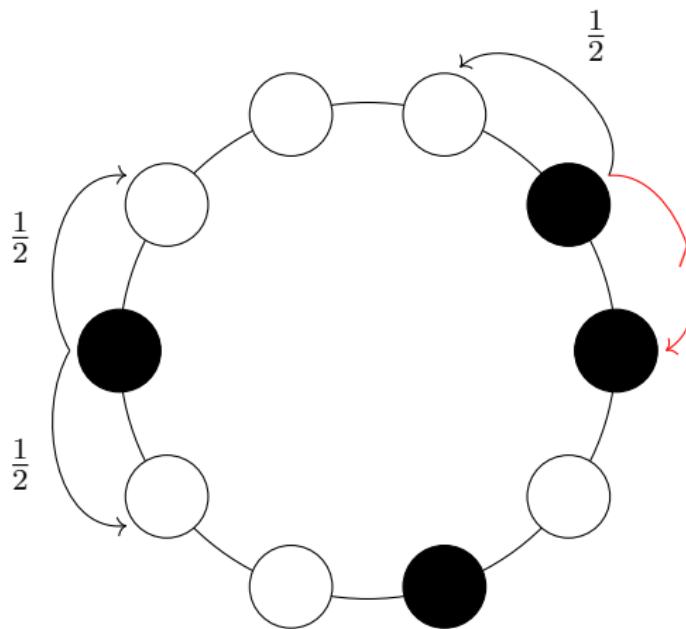
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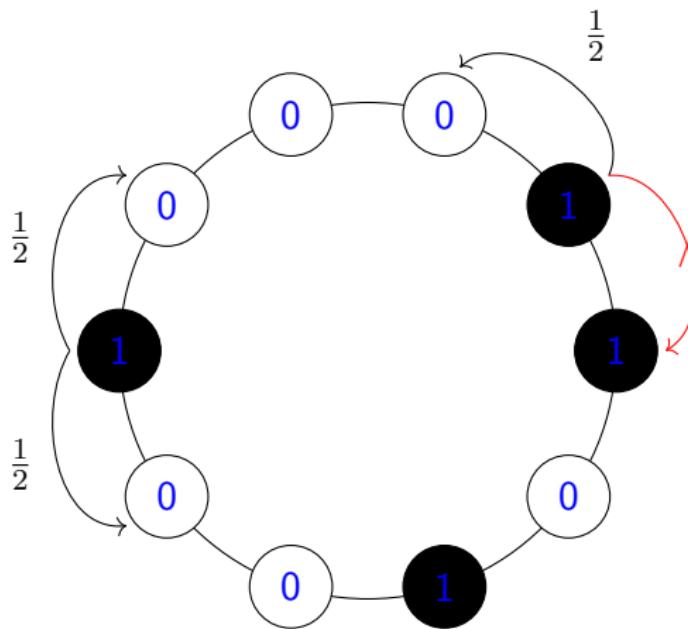
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# SSEP: Notation

- on torus  $\mathbb{T}^N := \mathbb{Z}/N\mathbb{Z}$
- state space: **configurations**  $\Omega^N := \{0, 1\}^{\mathbb{T}^N}$
- configuration  $\eta \in \Omega^N$ :

$\eta_x = 0 \Leftrightarrow$  site  $x$  is empty

$\eta_x = 1 \Leftrightarrow$  site  $x$  is occupied

- write  $\eta^{x,y}$  for the configuration where sites  $x$  and  $y$  are swapped:

$$\eta_z^{x,y} = \begin{cases} \eta_y & \text{if } z = x \\ \eta_x & \text{if } z = y \\ \eta_z & \text{otherwise} \end{cases}$$

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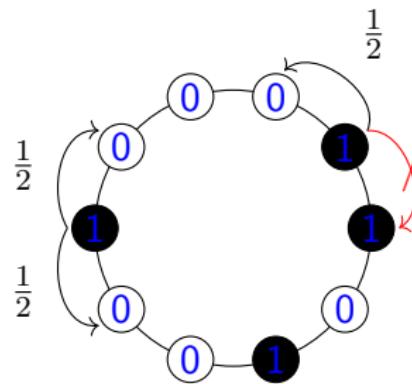
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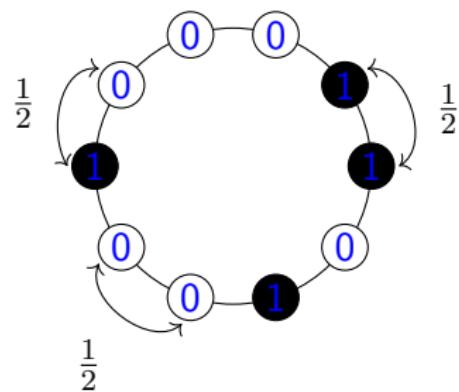
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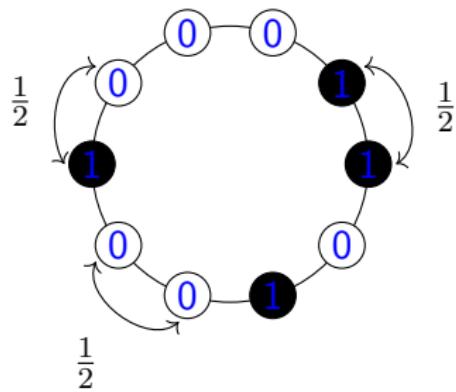
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# Two descriptions of the world

## A microscopic description

- atoms
- molecules
- cars
- microscopic particles

## A macroscopic description

- heat diffusion  
(heat equation)
- fluid dynamics  
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- traffic flow  
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# From microscopic to macroscopic

- zoom out:  $x \rightsquigarrow \frac{x}{N}$
- keep the global density of particles fixed!  $\rightsquigarrow \approx N$  particles
- accelerate time by  $\Theta(N)$
- take  $N \rightarrow +\infty$

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Instead of the mean density, use the **empirical distribution**:

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Models density:

$$\langle \pi^N(\eta), G \rangle = \frac{1}{N} \sum_{x=1}^N G\left(\frac{x}{N}\right) \eta_x \approx \int G(u) \rho^N(u) \, du$$

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LLN: Inside a **small macroscopic** ball around  $\frac{x}{N}$  are a lot of particles!

- ~ averaging effect
- ~ deterministic limit

Mixing: After mixing time, process is at **equilibrium**!

- ~ at time scale  $\Theta(N)$ : **locally at equilibrium**
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Three regimes:

- trivial: too small  $\rightsquigarrow$  no evolution
- hydrodynamic: exactly right  $\rightsquigarrow$  local equilibrium
- hydrostatic: too big  $\rightsquigarrow$  global equilibrium

LLN  $\leftrightarrow$  Mixing:

- Mixing happens in microscopic box of size  $\frac{\varepsilon \Theta(N)}{N}$   
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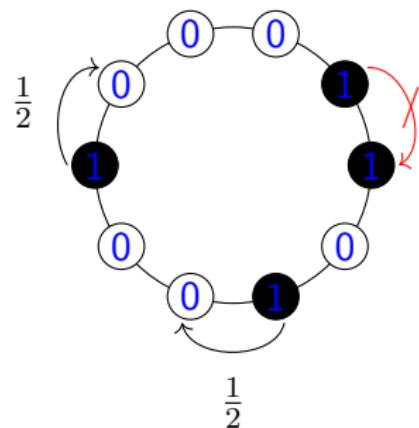
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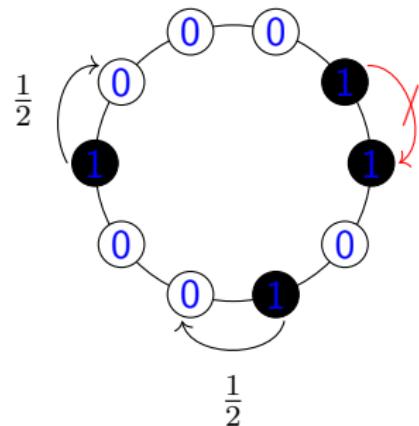
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$$\eta_x \eta_{x+1} \approx \overleftarrow{\eta}_x^\ell \cdot \overrightarrow{\eta}_{x+1}^\ell$$

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## References I



Claude Kipnis and Claudio Landim. *Scaling limits of interacting particle systems*. Grundlehren der mathematischen Wissenschaften. Springer, 1999.



Thomas M. Liggett. *Interacting Particle System*. 1st ed. Vol. 276. Classics in Mathematics. Springer, Berlin, Heidelberg, 1985.



Sunder Sethuraman and Doron Shahar.  
“Hydrodynamic limits for long-range asymmetric interacting particle systems”. In: *Electronic Journal of Probability* 23.none (2018), pp. 1–54. DOI: 10.1214/18-EJP237. URL: <https://doi.org/10.1214/18-EJP237>.

## References II



Lu Xu. *Hydrodynamic limit for asymmetric simple exclusion with accelerated boundaries*. 2021. arXiv: 2108.09345 [math.PR].



Lu Xu. *Hydrodynamics for one-dimensional ASEP in contact with a class of reservoirs*. 2022. DOI: 10.48550/ARXIV.2203.15091. URL: <https://arxiv.org/abs/2203.15091>.