Uniqueness of the Characteristic Initial Value Problem in General Relativity

Gabriel Sánchez-Pérez & Marc Mars

Department of Fundamental Physics University of Salamanca (Spain) Talk at Interdisciplinary junior scientist workshop





March 1, 2023

Junior scientist workshop

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Outlook

- Motivation
- 2 Hypersurface Data
- 3 Non-degenerate submanifolds
- 4 Double Null Data
- 5 Existence Theorem
- 6 Uniqueness Theorem
- Conclusions and future work

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• Then there exists a **Ricci-flat ambient spacetime** in which (Σ, h, K) is embedded as a spacelike hypersurface, with h and K being the first and second fundamental forms, respectively.

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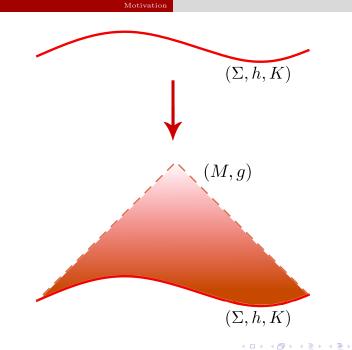
Motivation

 (Σ, h, K)

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ABSTRACT INITIAL DATA

EINSTEIN SPACETIME IN WHICH THE ABSTRACT DATA IS EMBEDDED

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Theorem

Isometric initial data give rise to isometric spacetimes.

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GEOMETRIC EXISTENCE THEOREM

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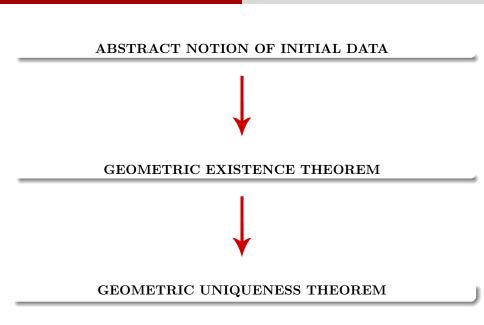
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A natural question

IS IT POSSIBLE TO DO THE SAME WITH THE CHARACTERISTIC CAUCHY PROBLEM IN GR?

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• The initial data is posed on two null and transverse hypersurfaces \mathcal{H} and $\underline{\mathcal{H}}$.

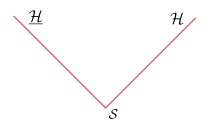
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- The initial data is posed on two null and transverse hypersurfaces \mathcal{H} and $\underline{\mathcal{H}}$.
- By providing $g_{\mu\nu}$ on $\mathcal{H} \cup \underline{\mathcal{H}}$ appropriately, $\exists (M, g)$ Ricci-flat in which the initial data is embedded.

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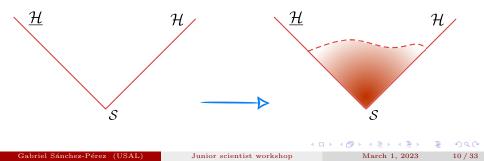
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"Standard" Cauchy Problem "Characteristic" Cauchy Problem

Abstract Initial Data

 (Σ, h, K)



(Abstract) Existence Theorem





(Abstract) Uniqueness Theorem





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A natural question

IS IT POSSIBLE TO DO THE SAME WITH THE CHARACTERISTIC CAUCHY PROBLEM IN GR?

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Answer

YES, BUT WE HAVE TO WORK A LITTLE BIT ...

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Define an abstract notion of a null hypersurface (null hypersurface data)

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- Define an abstract notion of a null hypersurface (null hypersurface data)
- ② Glue two null hypersurface data along the "corner" (double null data)

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YES, BUT WE HAVE TO WORK A LITTLE BIT...

- Define an abstract notion of a null hypersurface (null hypersurface data)
- ② Glue two null hypersurface data along the "corner" (double null data)
- **③** Prove **existence** and **uniqueness** theorems in this abstract framework

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Definition

• The Hypersurface Data formalism allows us to study any hypersurface of any causal character from an abstract point of view (see [M. Mars, 2012]).

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- A particular case is "Null Hypersurface Data". It consists of a set $\mathcal{D} = \{\mathcal{H}, \gamma, \ell, \ell^{(2)}, Y\}$ consisting of a smooth manifold \mathcal{H} , a 2-covariant, symmetric tensor field γ , a one-form ℓ , a scalar $\ell^{(2)}$ and a 2-covariant, symmetric tensor field Y, satisfying:

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- **2** Rad $(\gamma) \neq \emptyset$ and γ is semi-positive definite.
- This definition is **abstract**: it does not require the presence of any ambient spacetime.

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• This geometric construction is endowed with an internal **gauge freedom**.

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Gauge group

- This geometric construction is endowed with an internal gauge freedom.
- Given $\mathcal{D} = \{\mathcal{H}, \gamma, \ell, \ell^{(2)}, Y\}$ and $(z, \zeta) \in \mathcal{F}^{\star}(\mathcal{H}) \times \Gamma(T\mathcal{H})$ we define $\mathcal{D}' = \{\mathcal{H}, \gamma', \ell', \ell^{(2)'}, Y'\}$ by means of

$$\begin{split} \gamma' &\coloneqq \gamma, \\ \ell' &\coloneqq z \left(\ell + \gamma(\zeta, \cdot) \right), \\ \ell^{(2)\prime} &\coloneqq z^2 \left(\ell^{(2)} + 2\ell(\zeta) + \gamma(\zeta, \zeta) \right), \\ Y' &\coloneqq zY + \ell \otimes_s \mathrm{d}z + \frac{1}{2} \mathscr{L}_{z\zeta} \gamma. \end{split}$$

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- Let $\mathcal{D} = \{\mathcal{H}, \gamma, \ell, \ell^{(2)}, Y\}$ be NHD.
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- The tensors {γ, ℓ, ℓ⁽²⁾} are sufficient to reconstruct all the components of the spacetime metric on *H*.
- The tensor Y carries information of the derivative of g along a transverse direction.
- The gauge freedom is associated to the freedom in the choice of a transverse vector to the hypersurface (the so-called rigging vector ξ).

• It is possible to write down **all** tangential components of the ambient Ricci tensor in the embedded case in terms of the abstract data

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$$\operatorname{Ric}(g)_{ab} \stackrel{\mathcal{H}}{=} f_{ab}\left(\gamma, \ell, \ell^{(2)}, Y\right).$$

• The explicit form of f_{ab} is

$$f_{ab} = \left(\gamma_{af}\overline{R}^{f}{}_{cbd} + 2\overline{\nabla}_{[d}\left(K_{b]c}\ell_{a}\right) + 2\ell^{(2)}K_{c[b}K_{d]a}\right)P^{cd} - \left(\ell_{d}\overline{R}^{d}{}_{bac} + 2\ell^{(2)}\overline{\nabla}_{[c}K_{a]b} + K_{b[a}\overline{\nabla}_{c]}\ell^{(2)} + \ell_{d}\overline{R}^{d}{}_{abc} + 2\ell^{(2)}\overline{\nabla}_{[c}K_{b]a} + K_{a[b}\overline{\nabla}_{c]}\ell^{(2)}\right)n^{c}.$$

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• $K, P, n \text{ and } \overline{\nabla} \text{ are abstract (they are defined from } \{\gamma, \ell, \ell^{(2)}, Y\}).$

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$$\operatorname{Ric}(g)_{ab} \stackrel{\mathcal{H}}{=} f_{ab}\left(\gamma, \ell, \ell^{(2)}, Y\right).$$

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$$\mathcal{R}_{ab} \coloneqq f_{ab}\left(\gamma, \ell, \ell^{(2)}, Y\right).$$

- \mathcal{R}_{ab} makes no reference to any ambient space.
- This tensor allows us to define the constraint equations **abstractly** by

$$\mathcal{R}_{ab} = \lambda \gamma_{ab}$$

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Relation with Stefan's lecture

• In the appropriate gauge, and provided the data is embedded, one can recover the null structure equations

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$$\begin{aligned} \mathcal{R}(n,n) &= -n \left(\operatorname{tr}_{h} \chi \right) - 2\omega \operatorname{tr}_{h} \chi - |\chi|^{2}, \\ \mathcal{R}(n,X) &= -\left(\widetilde{\nabla}_{n} \tau \right)(X) + 2\nabla_{X}^{h} \omega - \left(\chi \cdot \tau \right)(X) - \tau(X) \operatorname{tr}_{h} \chi + \operatorname{div}(\chi)(X) \\ &- \nabla_{X}^{h} \operatorname{tr}_{h} \chi - 2\omega \nabla_{X}^{h} \log |\lambda| - \left(\chi \cdot d \left(\log |\lambda| \right) \right)(X), \\ \mathcal{R}_{AB} &= R_{AB}^{h} - 2\widetilde{\nabla}_{n} \Upsilon_{AB} + 4\omega \Upsilon_{AB} - \nabla_{A}^{h} \eta_{B} - \nabla_{B}^{h} \eta_{A} - 2\eta_{A} \eta_{B} \\ &- \chi_{AB} \operatorname{tr}_{h} \Upsilon - \Upsilon_{AB} \operatorname{tr}_{h} \chi, \\ H &= n \left(\operatorname{tr}_{h} \Upsilon \right) + |\eta|^{2} + \operatorname{div}(\eta) - \frac{1}{2} R^{h} + \left(\operatorname{tr}_{h} \chi - 2\omega \right) \operatorname{tr}_{h} \Upsilon, \end{aligned}$$

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Relation with Stefan's lecture

• In the appropriate gauge, and provided the data is embedded, one can recover the null structure equations

$$\begin{split} \mathcal{R}(n,n) &= -n \left(\operatorname{tr}_{h} \chi \right) - 2\omega \operatorname{tr}_{h} \chi - |\chi|^{2}, \\ \mathcal{R}(n,X) &= - \left(\widetilde{\nabla}_{n} \tau \right)(X) + 2 \nabla_{X}^{h} \omega - \left(\chi \cdot \tau \right)(X) - \tau(X) \operatorname{tr}_{h} \chi + \operatorname{div}(\chi)(X) \\ &- \nabla_{X}^{h} \operatorname{tr}_{h} \chi - 2\omega \nabla_{X}^{h} \log |\lambda| - \left(\chi \cdot d \left(\log |\lambda| \right) \right)(X), \\ \mathcal{R}_{AB} &= R_{AB}^{h} - 2 \widetilde{\nabla}_{n} \Upsilon_{AB} + 4\omega \Upsilon_{AB} - \nabla_{A}^{h} \eta_{B} - \nabla_{B}^{h} \eta_{A} - 2 \eta_{A} \eta_{B} \\ &- \chi_{AB} \operatorname{tr}_{h} \Upsilon - \Upsilon_{AB} \operatorname{tr}_{h} \chi, \\ H &= n \left(\operatorname{tr}_{h} \Upsilon \right) + |\eta|^{2} + \operatorname{div}(\eta) - \frac{1}{2} R^{h} + \left(\operatorname{tr}_{h} \chi - 2\omega \right) \operatorname{tr}_{h} \Upsilon, \end{split}$$

 \bullet Then, ${\cal R}$ generalizes them to the abstract setting and to any gauge.

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Given NHD D = {H, γ, ℓ, ℓ⁽²⁾, Y} and a non-degenerate submanifold S →
 H (i.e. i*γ ≕ h is a metric), we would like to define its two second fundamental forms and torsion one-forms.

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- We define a normal pair $\{C, T\} \in \mathcal{F}(\mathcal{H}) \times \mathfrak{X}(\mathcal{H})|_{\mathcal{S}}$ by means of $\gamma(T, X) + C\ell(X) = 0$ for all $X \in \mathfrak{X}(\mathcal{S})$.

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- When the data is embedded, $t \coloneqq \Phi_{\star}T + C\xi$ is **normal** to \mathcal{S} .

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- Given NHD D = {H, γ, ℓ, ℓ⁽²⁾, Y} and a non-degenerate submanifold S →
 H (i.e. i*γ =: h is a metric), we would like to define its two second fundamental forms and torsion one-forms.
- We define a normal pair $\{C, T\} \in \mathcal{F}(\mathcal{H}) \times \mathfrak{X}(\mathcal{H})|_{\mathcal{S}}$ by means of $\gamma(T, X) + C\ell(X) = 0$ for all $X \in \mathfrak{X}(\mathcal{S})$.
- When the data is embedded, $t \coloneqq \Phi_{\star}T + C\xi$ is **normal** to \mathcal{S} .
- Given a basis of NPs $\{t_1, t_2\}$ one can define two 2-covariant tensors \mathcal{K}^{t_1} , \mathcal{K}^{t_2} together with a one-form $\exists [t_1, t_2]$.

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$$\begin{aligned} \mathcal{K}^{\mathfrak{t}}(X,Z) &\coloneqq \frac{1}{2} \left(\mathscr{L}_{T} \gamma \right) (X,Z) + CY(X,Z) + \frac{1}{2} \left(X(C)\ell(Z) + Z(C)\ell(X) \right), \\ \overline{}_{ij}(X) &\coloneqq \left(\ell(T_{i}) + C_{i}\ell^{(2)} \right) \left(X(C_{j}) - K(X,T_{j}) \right) + C_{i}\ell \left(\overline{\nabla}_{X}T_{j} \right) \\ &+ \gamma(T_{i},\overline{\nabla}_{X}T_{j}) + \frac{1}{2}C_{i}C_{j}X(\ell^{(2)}) + C_{j}\Pi(X,T_{i}) \end{aligned}$$

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- When the data is embedded, $t := \Phi_{\star}T + C\xi$ is normal to S.
- \bullet Given a basis of NPs $\{\mathfrak{t}_1,\mathfrak{t}_2\}$ one can define two 2-covariant tensors $\mathcal{K}^{\mathfrak{t}_1},$ $\mathcal{K}^{\mathfrak{t}_2}$ together with a one-form $\exists [\mathfrak{t}_1, \mathfrak{t}_2].$
- When \mathcal{D} happens to be embedded, $\mathcal{K}^{\mathfrak{t}_1}$, $\mathcal{K}^{\mathfrak{t}_2}$ and $\exists [\mathfrak{t}_1, \mathfrak{t}_2]$ coincide with the two second fundamental forms and the torsion one-form in the basis $\{t_1 = T_1 + C_1\xi, t_2 = T_2 + C_2\xi\}.$

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$$\begin{aligned} \bullet & \phi^{\star}\underline{h} = h, \\ \bullet & \Psi^{\star}\left(\underline{\mathcal{K}}^{\Psi(\mathfrak{t}_{i})}\right) = \mathcal{K}^{\mathfrak{t}_{i}}, \\ \bullet & \Psi^{\star}\left(\underline{\neg}(\Psi(\mathfrak{t}_{i}), \Psi(\mathfrak{t}_{j}))\right) = \neg(\mathfrak{t}_{i}, \mathfrak{t}_{j}) \end{aligned}$$

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• They are **gauge-covariant** and independent on the basis of normal pairs.

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Existence theorem

Theorem (M. Mars, G. Sánchez-Pérez, arXiv:2205.15267)

Let $\{\mathcal{D}, \underline{\mathcal{D}}, \mu\}$ be double null data of dimension m > 1 satisfying the abstract constraint equations

$$\mathcal{R} = \frac{2\Lambda}{m-1}\gamma \quad and \quad \underline{\mathcal{R}} = \frac{2\Lambda}{m-1}\underline{\gamma},$$

where $\Lambda \in \mathbb{R}$. Then there exists a development (\mathcal{M}, g) of $\{\mathcal{D}, \underline{\mathcal{D}}, \mu\}$ (possibly restricted if necessary) solution of the Λ -vacuum Einstein equations. Moreover, for any two such developments (\mathcal{M}, g) and $(\widehat{\mathcal{M}}, \widehat{g})$, there exist neighbourhoods of $\mathcal{H} \cup \underline{\mathcal{H}}, \mathcal{U} \subseteq \mathcal{M}$ and $\widehat{\mathcal{U}} \subseteq \widehat{\mathcal{M}}$, and a diffeomorphism $\varphi : \mathcal{U} \longrightarrow \widehat{\mathcal{U}}$ such that $\varphi^* \widehat{g} = g$.

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Remember the spacelike case...

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Definition

 (Σ_1, h_1, K_1) and (Σ_2, h_2, K_2) were said to be isometric provided that $\exists \Sigma_1 \xrightarrow{\phi} \Sigma_2$ diffeo. such that $\phi^*\{h_2, K_2\} = \{h_1, K_1\}.$

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Now we have to take into account the **gauge freedom** too!

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Now we have to take into account the **gauge freedom** too!

Definition

We say that two double null data $\{\mathcal{D}, \underline{\mathcal{D}}, \mu\}$ and $\{\widehat{\mathcal{D}}, \widehat{\underline{\mathcal{D}}}, \widehat{\mu}\}$ are isometric if there exist diffeomorphisms $\psi : \mathcal{H} \longrightarrow \widehat{\mathcal{H}}$ and $\underline{\psi} : \underline{\mathcal{H}} \longrightarrow \underline{\widehat{\mathcal{H}}}$ and gauge parameters (z, ζ) and $(\underline{z}, \underline{\zeta})$ in \mathcal{D} and $\underline{\mathcal{D}}$, respectively, such that the pull-back double null data $\Xi^*\{\widehat{\mathcal{D}}, \underline{\widehat{\mathcal{D}}}, \widehat{\mu}\}$ satisfies

$$\Xi^{\star}\{\widehat{\mathcal{D}},\underline{\widehat{\mathcal{D}}},\widehat{\mu}\}=\mathcal{G}\left(\{\mathcal{D},\underline{\mathcal{D}},\mu\}\right).$$

Uniqueness theorem

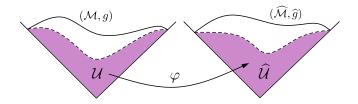
Theorem (M. Mars, G. Sánchez-Pérez, arXiv:2301.02722)

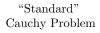
Let $\{\mathcal{D}, \underline{\mathcal{D}}, \mu\}$, $\{\widehat{\mathcal{D}}, \widehat{\underline{\mathcal{D}}}, \widehat{\mu}\}$ be two isometric DND satisfying the abstract constraint equations, and let (\mathcal{M}, g) , $(\widehat{\mathcal{M}}, \widehat{g})$ be respective developments. Then there exist neighbourhoods $\mathcal{U} \subseteq \mathcal{M}$ and $\widehat{\mathcal{U}} \subseteq \widehat{\mathcal{M}}$ of $\mathcal{H} \cup \underline{\mathcal{H}}$ and $\widehat{\mathcal{H}} \cup \underline{\widehat{\mathcal{H}}}$, respectively, and a diffeomorphism $\varphi : \mathcal{U} \longrightarrow \widehat{\mathcal{U}}$ such that $\varphi^* \widehat{g} = g$.

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"Characteristic" Cauchy Problem

Abstract Initial Data

 (Σ, h, K)

 $\{\mathcal{D}, \underline{\mathcal{D}}, \mu\}$

(Abstract) Existence Theorem





(Abstract) Uniqueness Theorem





Gabriel Sánchez-Pérez (USAL)

Junior scientist workshop

March 1, 2023

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- It puts the characteristic problem on the same footing as the standard Cauchy problem.
- It encompasses **all possible initial data** constructed from the metric and its first transverse derivative.
- It allows us to **work in any gauge** (not necessarily in the one in which the null structure equations are written).
- We want to study characteristic **Killing Initial Data** within this formalism (end-of-degree project with a student) and also the conformal field equations in arbitrary dimension.

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References

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