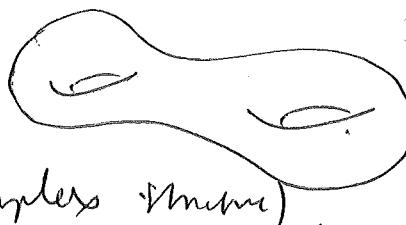


M surface genus \mathfrak{g} oriented



g Riem metric (\Rightarrow complex metric).

$$[g] = \{e^{\varphi}g; \varphi \in C^\infty(M)\} \text{ conf. class of } g (\Leftrightarrow \varphi \text{ shme})$$

(Gawedski) Def : a conformal field theory (CFT) are the following objects:

$Z: \text{Met}(M) \rightarrow \mathbb{R}^+$ partition function of the CFT.

$$(1) Z(e^{\varphi}g) = Z(g) \exp\left(\frac{G_Z}{96\pi} \int_M |d\varphi|^2_g + 2k_g \varphi \cdot d\omega_g\right).$$

k_g - scaling weight of g

G_Z - central charge of the theory.

and $\forall \varPhi \in \text{Diff}(M)$. $Z(\varPhi^*g) = Z(g)$.

(II) correlation function: $x_1, \dots, x_n \in M$, $\underbrace{\alpha_1, \dots, \alpha_n}_\text{weights} \in \mathbb{R}$.

$$Z(g; \underbrace{\varphi_{\alpha_1}(x_1), \dots, \varphi_{\alpha_n}(x_n)}_\text{primary fields}) \in \mathbb{R},$$

primary fields.

$$\text{Ideally: } Z(g) = \int_E e^{-S_g(\varPhi)} \underbrace{D\varPhi}_{\text{infinit & fund measure.}} \quad \begin{array}{l} | E \text{ infinite dim mfld.} \\ | E \cong \text{maps } (M \rightarrow \mathbb{R}). \end{array}$$

$S_g: E \rightarrow \mathbb{R}$ action

①

$$Z(g; \varphi_{x_1}(x_1), \dots, \varphi_{x_N}(x_N)) = \int_{\mathbb{H}^N} e^{-\sum_{i=1}^N \alpha_i \varphi(x_i)} e^{-S_g(\Phi)} D\Phi.$$

$$Z(e^{\varphi} g; \varphi_{x_1}(x_1), \dots, \varphi_{x_N}(x_N)) = Z(g; \varphi_{x_1}(x_1), \dots, \varphi_{x_N}(x_N)).$$

$$\times \exp\left(\frac{C\varepsilon}{96\pi} \cdot \int_M (|\partial\varphi|^2 g + 2h_g(\varphi)) dv_g\right).$$

$$\times \exp\left(\sum_{i=1}^N \Delta_i(\alpha_1, \dots, \alpha_N) \Psi(x_i)\right).$$

↓
eR, weights.

Renormalized volume of 3 dim hyperbolic wfs : $M = M^{3,2}$.

$$\text{Gauss-Bonnet } \int_M h_g \cdot dv_g = 4\pi(2 - 2g).$$

Conf. class for h_g : if $\hat{g} = e^w g$.

$$h_{\hat{g}} = e^{-w} (\Delta_g w + h_g).$$

Thm. $\forall g$ metric on M , $\exists! w \in C^\infty(M)$ s.t.

$\hat{g} = e^w g$ has scalar curv. $K_{\hat{g}} = -2$.

(M, \hat{g}) can be realized as $\mathbb{H}^2 \backslash \mathbb{H}^2$ half plane,

$\Gamma \subset \mathrm{PSL}_2(\mathbb{R})$, discrete subgroup.

Pf. use variational method:

let $f_g: C^\infty(M) \rightarrow \mathbb{R}$,

$$\Psi \mapsto \int_M \left(\frac{1}{2} |\partial \Psi|^2 + k_g \Psi + 2e^{\Psi} \right) dv_g .$$

minimise Ψf_g , find Ψ_0 , $\hat{g} = e^{\Psi_0} g$ has $k_{\hat{g}} = -2$.

Ψ_0 - unique minimiser, critical pt.

Remark. if we fix $\text{vol}(M, e^{\Psi_0} g) = 1$. $\int_M e^{\Psi_0} dv_{\hat{g}} = 1$, you can get Ψ_0 by minimising $\int \left(\frac{1}{2} |\partial \Psi|^2 + k_g \Psi \right) dv_g$.

\Rightarrow partition func. of a CFT is minimised at the metric with such value $= -2$. so fun

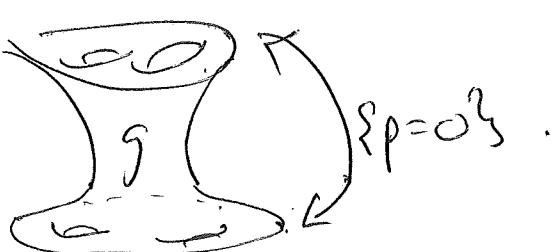
$$\int u \log dv_g = \int_{-2}^{GB} dv_g \stackrel{?}{=} 4\pi(1-2G).$$

Now let (X, g) be a 3-dim hyp. mfld non-cpt but compactly compact. I.e. $\exists \bar{X}$ mfld with holes s.t. $\partial \bar{X} = M$ compact, and with \bar{X} cpt.

$\exists \rho: C^\infty(\bar{X}) \rightarrow \mathbb{R}^+$ bdy sloping function.

$\{p=0\} = \partial \bar{X}$, $p > 0$ in X , $d\rho$ does not vanish.

on $\partial \bar{X}$



$$\overset{\circ}{X} = X.$$

③

S.t. P^2g is a metric on \bar{X} .

Ques. a short computation using $\text{curl}(g) = -1$ near $\partial\bar{X}$.

$[dg]_{P^2g} = -1$ at $\partial\bar{X}$. (Assume X compact).

X can be realized as $T\mathbb{H}^3$, $T \subset \text{Isom}(\mathbb{H}^3) = \text{PSL}_2(q)$.

$\text{PSL}_2(q)$ acts on $\mathbb{B}^3 = \{m \in \mathbb{R}^3 \mid m| \leq 1\}$ as conformal transform. Let $\Lambda(T) = \overline{\{r(x) : x \in T\}}^{\mathbb{B}^3}$ linear hull of $T \subset \mathbb{S}^2 = \partial\mathbb{H}^3$.

Let $x \in \mathbb{H}^3$; $r(T) = \mathbb{S}^2 \setminus \Lambda(T)$ qm, T acts properly discontinuously.

$\bar{X} = T \setminus (\mathbb{B}^3 \setminus \Lambda(T))$. proper & discontinuously.

$\partial\bar{X} = T \setminus r(T) = \bigsqcup_{i=1}^N M_i$ \leftarrow connected Riemann surfaces. It has the complex confor. st. of $T \setminus r(T)$.

M_i has a projective structure:

i.e. an atlas ~~of~~ with charts $\varphi_i : M_i \rightarrow \mathbb{CP}^1$,

$\varphi_n \circ \varphi_j^{-1}$ is in $\text{PSL}_2(q)$.

(Let confnd charts on $M = \bigsqcup_{i=1}^N M_i = \partial\bar{X}$,

$\underbrace{[P^2g]_{T_M}}_{\bar{g}}$. Confnd boundary.

Lemma (Graham) $\forall h_0 \in [g^2 g]_{T\partial M}$, $\exists!$ non $\bar{\partial}X$ bdy defin from $\hat{p} = e^w p$ s.t. $|d\hat{p}|_{\hat{p}^2 g} = 1$. non $\bar{\partial}X$ and $\hat{p}^2 g|_{T\partial M} = h_0$.

Consequence: if we take the flow of $\nabla_{\hat{p}}^{\hat{g}^2}$; we get a diffn: $\Psi: (0, \epsilon)_{\hat{p}} \times \bar{\partial}X \rightarrow M \subset \bar{\partial}X$.

$$\Psi^* g = \frac{dp^2 + h(\hat{p})}{\hat{p}^2}, \quad h \text{ is a family of metrics on } \bar{\partial}X.$$

Lemma (Teuffel-Braun). $h(\hat{p}) = h_0 + \hat{p}^2 h_2 + \hat{p}^4 h_4$.

h_2 forms on $\bar{\partial}X$. with $\text{Tr}_{h_0}(h_2) = -\frac{1}{2} k_{h_0}$.

$$\text{and } S_{h_0}(h_2) = \frac{1}{2} \cdot dK_{h_0}.$$

$$h_4 = \frac{1}{4} h_2^2 \quad (h_2 \text{ defin by annih } h_2 \text{ w.r.t. end})$$

h_2 fully under determined, but trace is determined.

Lemma (Graham) - Recall: (x, g) uniformly opt, hyperbolic 3mfld, ρ holo defining function, $\rho^2 g = \bar{g} \in C^\infty(\bar{x})$ in cech form

$$[\bar{g}|_{\partial \bar{x}}] \rightarrow \hat{h}_0, \exists! \text{ (new holo) } \hat{\rho} \text{ b.d.f. s.t. } |\hat{d\rho}|_{\bar{g}} = 1 \text{ near } \bar{x}.$$

At. let $\hat{\rho} = e^w \rho$ $w \in C^\infty(\bar{x})$, and $|\hat{d\rho}|_{\bar{g}} = 1$.

$$\Leftrightarrow 1 = \underbrace{\bar{g}(d\rho, d\rho)}_{1 + PQ, \text{ smooth.}} + \rho^2 |\hat{dw}|_{\bar{g}}^2 + 2\bar{g}(d\rho, dw)\rho.$$

$$\Leftrightarrow \text{ and } w|_{\partial \bar{x}}, \text{ then } e^{2w|_{\partial \bar{x}}} h_0 = \hat{h}_0 \\ e^{2w_0} h_0 = \hat{h}_0.$$

$$\Leftrightarrow 2(\nabla^{\bar{g}})(w) + Q + \rho |\hat{dw}|_{\bar{g}}^2 = 0.$$

Non-degenerate Hamilton-Jacobi eqⁿ.

Use grad-flow of $\hat{\rho}$ w.r.t. $\hat{\rho}^2 g$; get diffeomorphism.

$$\Psi: [t, \infty) \times \partial \bar{x} \rightarrow M \subset \bar{x}, \text{ s.t. }$$

$$\Psi^* g = \frac{d\rho}{\hat{\rho}^2} + \frac{\hat{h}^*(\hat{\rho})}{\hat{\rho}^2} \text{ moves on } \partial \bar{x}.$$

$\hat{\rho}$ is called as geodetic bending defining func. for (x, g) associated to \hat{h}_0 .

Now let ρ and $\hat{\rho}$ are no the geodetic b.d.f, s.t.

$$\hat{\rho} = e^w \rho, \quad \nabla_{\hat{\rho}}^{\bar{g}} = \partial_{\hat{\rho}}.$$

$$2\partial_p w = -\rho \left(1 \partial_p w^2 + |dy_0 w|_{h(\rho)}^2 \right), \quad \text{wh } |w|_{p=0} = w_0.$$

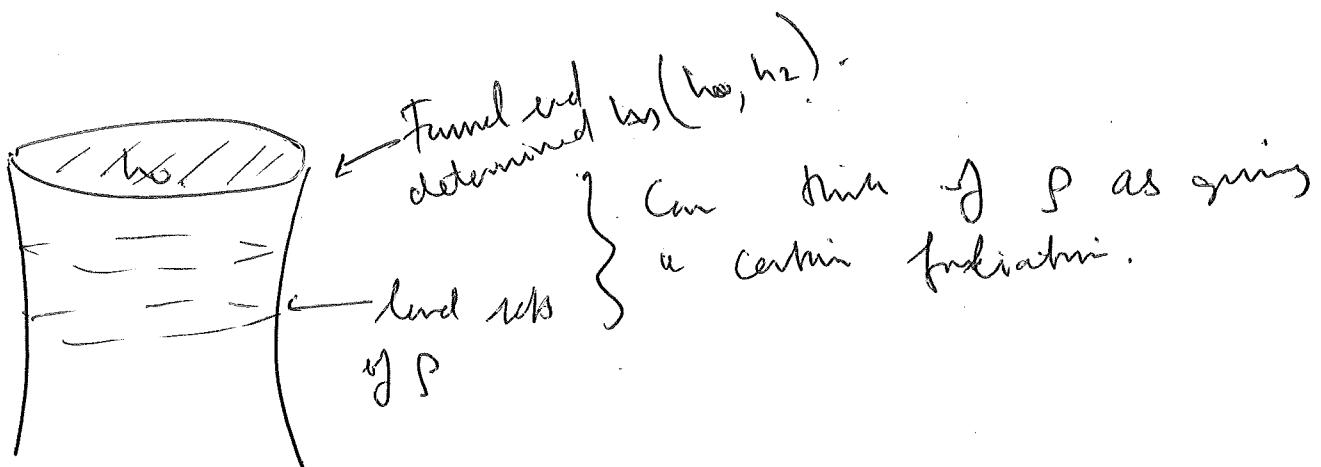
y coordinates on $\partial\bar{\mathcal{X}}$.

Taylor of w :

$$\text{(at } p=0) \quad w(p, y) = w_0(y) + p w_1(y) + p^2 w_2(y) + O(p^3).$$

$$\Rightarrow \text{get } 4p w_2 + O(p^3) = -\rho \left(4\rho^2 w_2^2 + |dy_0 w_0|_{h_0}^2 + 2\rho^2 \langle dw_0, dw_2 \rangle_{h_0} \right) + O(p^3)$$

$$\Rightarrow w_2 = -\frac{1}{4} |dw_0|_{h_0}^2 \text{ and } w_1 = 0.$$



If p half associated to $h_0 \in [p^2 g]_{T\bar{\mathcal{X}}}$, then:

$$h(p) = h_0 + p^2 h_2 + p^4 h_4.$$

$$h_4 = \frac{h_2 \circ h_2}{4} \quad (h_2 \text{ is damped with End } (T\bar{\mathcal{X}})).$$

and: $\begin{cases} \text{Tr}_{h_0}(h_2) = -\frac{1}{2} k_{h_0}, \\ \delta_{h_0}(h_2) = \frac{1}{2} d k_{h_0}. \end{cases}$ ← scalar curv.

$$\begin{cases} \text{Tr}_{h_0}(h_2) = -\frac{1}{2} k_{h_0}, \\ \delta_{h_0}(h_2) = \frac{1}{2} d k_{h_0}. \end{cases}$$

h_2 non-local in h_0 when Schrödinger Einstein eq. in X with hole condition [h_0].

choose h_0 with $k_{h_0} = 2 \Rightarrow \text{Tr}_{h_0}(h_2) = -1$ and $\delta_{h_0}(h_2) = 0$.

(2)

Teichmüller space of M h^{rt} is surface, trace free form (for h_0). in $T(M) = \text{Met}_{-1}(M) / \text{Diff}_0(M)$.

hypothetic .

diffs homotop
to identity.

$$\dim_{\mathbb{R}} T(M) = 6g - 6 + \text{q-mfd}, \quad T^* T(M) = \left\{ K \in \text{sym. 2-tens.} \right. \\ \left. \text{s.t. } \text{div}_{h_0}(K) = 0 = T_{h_0}(K) \right\}$$

$$\text{ht of form ends} / \text{Diff}_0(M) \cong T^* T(M).$$

Renormalized volume: the vol form of g near ∂X is

$$d\text{vol}_g = \frac{dp \cdot d\text{vol}_{h_0}(p)}{p^3} = \frac{dp}{p^3} d\text{vol}_{h_0} \times N(p).$$

$$N(p) = \det(\text{ch}_0^{-1} h(p)) = 1 + p^2 N_2 + O(p^4)$$

$$N_2 = \text{Tr}_{h_0}(h_0) = -\frac{1}{2} k_{h_0}.$$

Def. (Renormalized vol) (X, g) hyp. 3-mfd, h_0 choice of metric at ∂X .

$$\text{Vol}_R(X, g, h_0) = \underbrace{\text{FP}_{\varepsilon \rightarrow 0}}_{\text{finite part}} \left[\int_{\mathbb{R}^2} p^2 d\text{vol}_g \right].$$

Finite part, removing residue.

$$\underline{\text{Ex.}} \quad = \text{FP}_{\varepsilon \rightarrow 0} \int_{p > \varepsilon} p^2 d\text{vol}_g.$$

$$\underbrace{\quad}_{C_2 \varepsilon^{-2} + C_1 \log(\varepsilon) + \text{vol}_R(X, g, h_0)} + O(\varepsilon). \quad (3)$$

$$\text{Lima} \quad \text{vol}_R(X, g, e^{w_0} h_0) = \text{vol}(X, g, h_0) - \frac{1}{8} \int \left(\frac{(dw_0)^2}{2} + k_{h_0} w_0 \right) dw_0$$

"Conformal anomaly of a CFT."

$$\begin{aligned} \text{Hint of pf: } \text{FP} \int_{z=0} \hat{p}^2 dw_g &= \text{FP} \int_{z=0} \hat{p}^2 e^{zw} v(p) dp dw_{h_0} \\ &\quad + zw + O(z^2) \leftarrow \\ &= \text{vol}_R(X, g, h_0) + \underbrace{\text{res} \int_{z=0} \hat{p}^{2-3} N(p) w dw_{h_0} dp}_{\text{res}} \end{aligned}$$

$$\cancel{\text{res}} \int_{z=0} \hat{p}^2 (w_2 + w_0 N_2) \frac{dp}{z} dw_{h_0} = \cancel{\text{res}} \int_{\partial X} (w_2 + w_0 N_2) dw_{h_0}.$$

If $Z(g)$ is partition function of CFT, define the tensor.

$$Hg \in C^\infty(M, \otimes^2 T^*M), \quad ?$$

$$\langle d(\log Z(g)), h \rangle = \int_M \langle Hg, h \rangle_g dv_g.$$

↓
sign 2-form

$$Z(e^{wg}) = Z_g e^{\frac{c}{96\pi}} \int_M (1 dw_g^2 + 2 kg w) dw_g \quad \begin{array}{l} (2) \\ Z(T^*g) = Z(g) \\ \forall g \in \text{Diff}(M). \end{array}$$

①.

$$② \Rightarrow \delta_g(Hg) = 0 \text{ since } \langle Hg, L_x g \rangle = 0 \quad \forall x \text{ v. fields.}$$

$$① \Rightarrow T_{rg}(Hg) = \frac{2c}{96\pi} kg.$$

Call Hg stress-energy tensor.

(4)

If g_t is a family of hyp metrics on M ,

$$\partial_t|_{t=0} \mathcal{I}(g_t) = H_{g_0} \Rightarrow H_{g_0}^{\text{tr}} \text{ is div-free.}$$

$$\text{Tr-vec } H_{g_0}^{\text{tr}} \in T_{g_0}^* C(M).$$

Ex $\text{vol}_R(x, g, E^{**} h_0) = \text{vol}_R(x, g, h_0)$.

$\forall \mathbb{E}$ homotopic to identity.

Define $\text{vol}_R(x, g) = \text{vol}_R(x, g, h_0)$. hypothetic in conformal
holo.

Th If g_t is a family of conformally compact hyp 3mfld,

$$\partial_t|_{t=0} \text{vol}_R(x, g_t) = -\frac{1}{4} \int_{\partial\bar{x}} \langle h_0^{\text{tr}}, h_0 \rangle_{h_0} \text{div}_x,$$

$h_0 = \partial_t|_{t=0} h_0^t$. h_0^t hyp metric at $\partial\bar{x}$ for g_t .

PF. Use Schläfli formula: $\partial_t \text{vol}_R(x, g_t) = \frac{1}{n} \int_{\partial\bar{x}} (H + \frac{1}{2} \langle g, II \rangle_g) \text{div}_x$

If g_t is Einstein on cpt manifold with holo ($\dim n+1$)
t domain. ↑ 2nd ff.
of mean curv. at $\partial\bar{x}$.

* Can view vol_R as a function in Teichmller space.
over boundary;

Th: $\bar{\partial} \partial \text{vol}_R(x) = \mathcal{G}_{WP} \leftarrow$ Weil-Petersson. Symplectic fm.

Z partition function of 2D CFT $\text{Metrics } (M) \rightarrow \mathbb{R}$.

- Conformal anomalies, diffeo invariance.

$\rightarrow Z$ minimised at $g = g_{\text{hyp}}$ hyperbolic metric.

$dZ(g_{\text{hyp}})$. in the direction of Teichmüller space $T(M) = \text{Met}_+(M)$

$\stackrel{\text{II}}{\uparrow} \quad \frac{\delta Z}{\delta g_{\text{hyp}}} \quad , \quad A_{\text{hyp}}^{\text{tot}}$ is traceless & div-free, $\in T_{g_{\text{hyp}}}^* T(M)$.

" Stress-Energy".

Also, $T^* T(M) \xrightarrow{\text{iso}} P(M) = \text{set of projective symms on } M$.

\uparrow
from $Z \rightarrow$ a projective symm. section of
 $P(M) \rightarrow T(M)$ (affine bundle).

(See away has dual duality on Riem. metric)

(II). Liouville QFT.

(M, g) . Riem. surface, Liouville action:

$$S_L(g, \varphi) := \frac{1}{4\pi} \int \left(16\pi R_g + Q k_g \varphi + 4\pi \mu e^{2\varphi} \right) dv_g.$$

$S_L(g, \cdot)$ is minimised at φ_0 s.t. $g' = e^{2\varphi_0} g$ has

constant scalar curvature. if $Q = \frac{2}{\gamma}$, $k_g = -2\pi\mu\gamma^2$.

①

Liouville QFT:

$$Z(g) := \int_E e^{-S_L(g, \psi)} d\psi \quad \text{Formal integral.}$$

$\{\text{maps } m \rightarrow \mathbb{R}\}$

More generally: $Z(g, F) := \int F(\psi) e^{-S_L(g, \psi)} d\psi$, $F: E \rightarrow \mathbb{R}$,
in particular, with $F_i(\psi) = \prod_{i=1}^N e^{\alpha_i \psi(x_i)}$. primary fields.

Δ Rem: in physics, α should be $\frac{1}{2} + \frac{1}{2}$.

1) Gaussian Free Field: (cf. survey paper by Scott Sheffield).

$(\psi_i)_{i \in \mathbb{N}}$ o.n. basis of eigenfunctions of Δ_g , Δ_g

$$\Delta_g \psi_i = \lambda_i \psi_i, \quad \lambda_0 = 0, \quad \psi_0 = \frac{1}{\sqrt{\text{vol}_g(M)}}.$$

$(\Omega, \mathcal{F}, \mathbb{P})$ probability space,

$(\alpha_i)_{i \in \mathbb{N}}$ family of Gaussian random variables which
identically distributed independent $\mathbb{E}(\alpha_i) = 0$, $\mathbb{E}(\alpha_i, \alpha_j) = \delta_{ij}$.

Def GFF: $X_g(n) = \sum_{i=1}^{\infty} \alpha_i(n) \cdot \frac{\psi_i(n)}{\sqrt{\lambda_i}}$.

Rem: $X_g \in H^s(M)$ $\forall s > 0$ a.s.

$$\mathbb{E} \left(\| (1 + \Delta_g)^{-s/2} X_g \|_2 \right) = \mathbb{E} \left(\sum_{i=0}^{\infty} \frac{\alpha_i}{((1 + \lambda_i)^{s/2} \lambda_i^{1/2})} \right) = \sum_{i=0}^{\infty} \dots$$

By Weyl law. (2)

Covariance of X_g : $\varphi, \varphi' \in C^\infty(M)$.

$$\mathbb{E}(\langle X_{g,\varphi} \rangle \langle X_{g,\varphi'} \rangle) = \dots = 2\pi \langle G_g, \varphi \otimes \varphi' \rangle_{H^0(M)}.$$

where. $G_g(x, x') = \sum_{i \geq 1} \frac{\varphi_i(x)\varphi_i(x')}{\lambda_i}$

Kernell for op R_g $\Delta_g R_g = \mathbb{I}_d - \Pi_0 \leftarrow$ project on $N(\Pi_0)$.

Rem. $G_g(x, x') = \underbrace{\frac{1}{2\pi} \log(\text{dist}(x, x'))}_{\text{Rindist}} + m_g(x) + o(1).$ \uparrow
 "mass of g ".

Let $H_0^{-s}(M) = \{ \varphi \in H^s(M) : \langle \varphi, 1 \rangle_g = 0 \}$.

Lemma. \exists probability P on $H_0^{-s}(M)$ s.t. the law of random var X_g in P and $\forall \varphi \in H^s(M)$, $\langle X_g, \varphi \rangle$ is a random ^(gaussian) var on \mathbb{R} with zero mean and covariance $2\pi \langle G_g, \varphi \otimes \varphi \rangle$.

$$U \subset H_0^{-s}(M), P(X_g \in U) = P(U).$$

P represents the formal measure $F \mapsto \int F(\varphi) e^{-\frac{1}{4\pi} \int_M \|d\varphi\|^2} d\varphi$
 $\times \sqrt{\det(\frac{1}{2\pi} \Delta_g)^s}$.

Approximate on $1 \leq i \leq N$:

$$\int F(\varphi) e^{-\frac{1}{4\pi} \|d\varphi\|^2} d\varphi = \int_{\mathbb{R}^N} F\left(\sum_{1 \leq i \leq N} \alpha_i \varphi_i\right) \prod_{i=1}^N e^{-\frac{1}{4\pi} \frac{\alpha_i^2}{\lambda_i} \lambda_i} d\alpha_i.$$

$\delta > 0$ fixed. To get a measure on $H^{-S}(M)$, take

$$\varphi: \begin{array}{ccc} H_0^{-S}(M) \times \mathbb{R} & \xrightarrow{\text{isom}} & H^{-S}(M) \\ & \longleftarrow & \text{indep. of } g \text{ as a.} \\ (x, c) & \mapsto & x + c. \end{array}$$

vec. space.
depends on g !

Def. let $P' = \varphi_{**}(P \otimes dc)$ ^{heuristic}

Lemma: P' in $H^{-S}(M)$ is conformally invariant:

$$\text{if } \cdot \hat{g} = e^w g; \int_M \mathbb{E}(F(x_{\hat{g}} + c)) dc = \int_M \mathbb{E}(F(x + c)) dc.$$

pf. use relation b/w G_g and $G_{\hat{g}}$.

$$P' \text{ represents: } \sqrt{\det(\frac{1}{2\pi} \Delta_g)} \cdot e^{-\frac{\|d\varphi\|_{L^2}^2}{2}} D\varphi.$$

wanted, $\varphi \in H^S$, by using fs.
by some magic.

is a measure on $H^{-S}(M)$.

To define LQFT: partition function

$$T_{q,r}(g) = \left(\frac{\det \Delta_g}{\text{vol}_g M} \right)^{\frac{1}{2}} \int_M \mathbb{E} \left(e^{\frac{-Q}{4\pi} \sum \log(x_{\hat{g}} + c) - r(c + x_g)} \right) d\mu_g$$

does not
matter since.

$$\text{with } Q = \frac{r}{r} + \frac{r}{2}.$$

makes sense since
 $x_g \in C^\infty$.

go through renormalisation. since $e^{-r(c + x_g)}$
multi. br.
distribution.

(4)

2) Gaussian multipl. came. obsns. (GMC).

(Kahane 85, Polter-Verguts 2008, Duplantier-Sheffield '09).

Let $X_{g,\varepsilon}(z) := \int X_g \cdot d\mu_{z,\varepsilon}$.
uniform mass on the
geodesic circle of radius
 $\varepsilon > 0$: centre $z \in \mathbb{H}$ for g .
Off in ∞ a.s.

Prop. If $r \in (0,2)$, then the measur.

$$g_y^r(z) = e^{r^2/2} e^{2X_{g,\varepsilon}(z)} dv_g(z).$$

Converges to a measure. (in prob & weak sense).

g_y^r called GMC.

Defin part. func. meas $M e^{R(-*X_g)}$ replaced by $M e^{R(g_y^r)}$.

Th. $\Pi_{\mu,r}$ is the partition function of a CFT meas.

current density $1 + 6Q^2$ if $Q = \frac{\partial}{2} + \frac{2}{\pi}$

& can define current function in the same way.